



ME 440

Intermediate Vibrations

Tu, March 12, 2009

Chapter 4: Vibration Under General Forcing Conditions

Section 4.4



Before we get started...

- Last Time:

- Periodic Excitation
- Special Excitation Cases

- Today:

- HW Assigned (due March 26):

- 4.23 (Use Eq. 4.36)
- 4.24 (See Example 4.9 in the book, it helps...)

- Material Covered:

- Response under an impulsive force
- Duhamel's Integral (convolution integral)

[New Topic: Case C (see slide beginning of lecture)]

Impulse Excitation and Response

- Impulse Excitation: What is it?
 - When a dynamical system is excited by a suddenly applied and short in duration nonperiodic excitation $F(t)$
 - Specifically, what does “short in duration” mean?
 - It means that it's shorter than 1/10 of the natural period ($\tau_n = 2\pi/\omega_n$)
 - This is a rule of thumb...
 - Note: steady state response not produced (forcing term is not periodic)
 - Only transient response present in the system
 - Oscillations take place at the natural and/or damped natural frequencies of the system

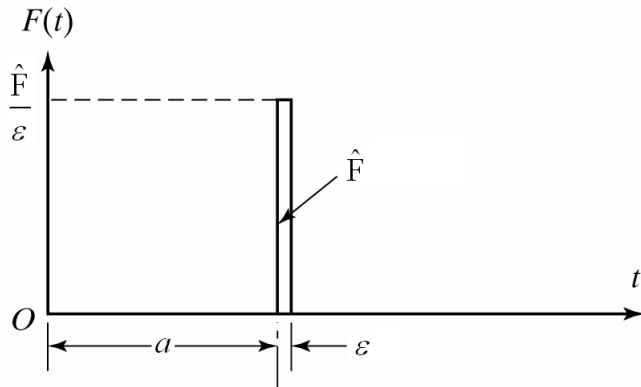
[Cntd.]

Impulse Excitation and Response

- Impulse: Time integral of the force

$$\hat{F} = \int F(t) dt$$

- Impulsive Force: A force of very large magnitude which acts for a very short time and has a finite *nonzero* impulse



As $\epsilon \rightarrow 0$, $\frac{\hat{F}}{\epsilon} \rightarrow \infty$,

but \hat{F} is finite.

- Impulse is the area of the rectangle in the figure above
 - "Spike" can be high, but area underneath is finite



Dirac Delta Function

- The type of impulsive function $F(t)$ for which the following condition holds:

$$\hat{F} = \int F(t)dt = 1$$

- Dirac Delta Function also called unit impulse, or delta function
- A delta function at $t=a$ is identified by the symbol $\delta(t-a)$ and has the following properties

$$(I) \quad \delta(t - a) = 0 \quad \text{for all } t \neq a$$

$$(II) \quad \int_0^{\infty} \delta(t - a)dt = 1 \quad 0 < a < \infty$$

Comment on Dirac Function

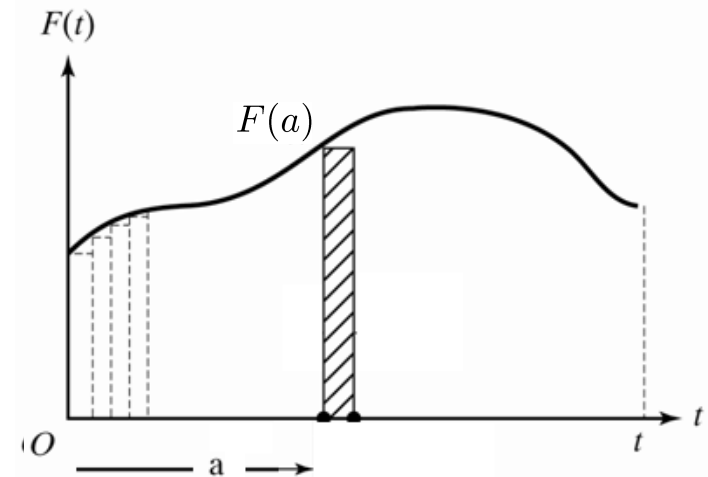
(somewhat tricky...)

- If you take a function $F(t)$ and multiply it by the Dirac Function $\delta(t-a)$, what do you get?

$$G(t) = F(t) \cdot \delta(t - a) \quad \text{for } \forall a \text{ with } 0 < a < \infty$$

- You get a function $G(t)$, which has two interesting properties:
 - $G(t)$ is zero everywhere, except at $t=a$
 - For any $0 < a < 1$, its integral from zero to 1 assumes the value $F(a)$:

$$\int_0^{\infty} G(t) dt = \int_0^{\infty} F(t) \cdot \delta(t - a) dt = F(a) \quad \text{for } 0 < a < \infty$$





Going Back To Mechanical Systems...

- Recall Newton's Second Law: $F(t) dt = m dv$

- Then,
$$\int_0^T F(t) dt = m \int_0^T dv \quad \Rightarrow \quad v(T) = v(0) + \frac{\hat{F}}{m}$$

- Assume now that $F(t)$ is an impulsive force applied at time $t=0$ with the goal of "kicking" the system in motion. Then,

$$v(0^+) = v(0) + \frac{\hat{F}}{m}$$

- If $v(0)=0$, then $v(0^+) = \frac{\hat{F}}{m}$

- Important observation: an impulsive force will lead to a sudden change in velocity without appreciable change in its displacement



Response To Impulsive Force ~ Undamped Systems ~

- Impulsive force applied to undamped system, free* response
- In Chapter 2, we saw that

$$x_h(t) = x_0 \cos \omega_n t + \frac{\dot{x}_0}{\omega_n} \sin \omega_n t$$

- Assume zero initial conditions, but system is kicked in motion by an impulsive force. Then,

$$x_0 = 0 \quad \dot{x}_0 = \frac{\hat{F}}{m}$$

- Response ends up looking like this:

$$x_h(t) = \frac{\hat{F}}{m\omega_n} \sin \omega_n t \quad \text{where} \quad \omega_n = \sqrt{\frac{k}{m}}$$



Response To Impulsive Force ~ Underdamped Systems ~

- Impulsive force applied to underdamped system, free* response
- In Chapter 2, we saw that

$$x_h(t) = e^{-\zeta\omega_n t} (B_1 \cos \omega_d t + B_2 \sin \omega_d t)$$

- Assume zero initial conditions, but system is kicked in motion by an impulsive force. Then,

$$x_0 = 0 \quad \dot{x}_0 = \frac{\hat{F}}{m}$$

- Response ends up looking like this:

$$x(t) = \frac{\hat{F}}{m\omega_d} e^{-\zeta\omega_n t} \sin \omega_d t$$



Follow Up Discussion, Concluding Remark

- Important observation, which simplifies the picture
 - Because we are dealing with linear systems, to find their response to any impulsive force, first find the response of the system to the Dirac Function
 - Response of the system to the Dirac Function is denoted by $g(t)$
 - Then, for any impulsive function with \hat{F} , the response is simply going to be a scaling of the response obtained for the Dirac Function:

$$x(t) = \hat{F} g(t)$$

- Look back at the previous two slides to see that the response of an underdamped or undamped system is indeed obtained as indicated above...

[New Topic:]

Arbitrary Excitation

- To understand this topic, you need to keep in mind three things:
 - First, we are dealing with linear systems, and the principle of superposition holds
 - Second, recall that an integral is defined as a sum (a special sum, that is)
 - Third, recall that $g(t)$ represents the response of the system to a Dirac Delta function. For any arbitrary impulse, the response is just a scaling of $g(t)$ to obtain $x(t)$
- The claim is that having the expression of $g(t)$ is a very important step in understanding the response of the system to *any* arbitrary external excitation $F(t)$

[Cntd.]

Arbitrary Excitation

- The fundamental idea:

- Consider the arbitrary excitation $F(t)$ to be just a train of successive impulsive excitations...
- ... then apply the principle of superposition to find the response as the sum of the responses to this train of impulsive excitations

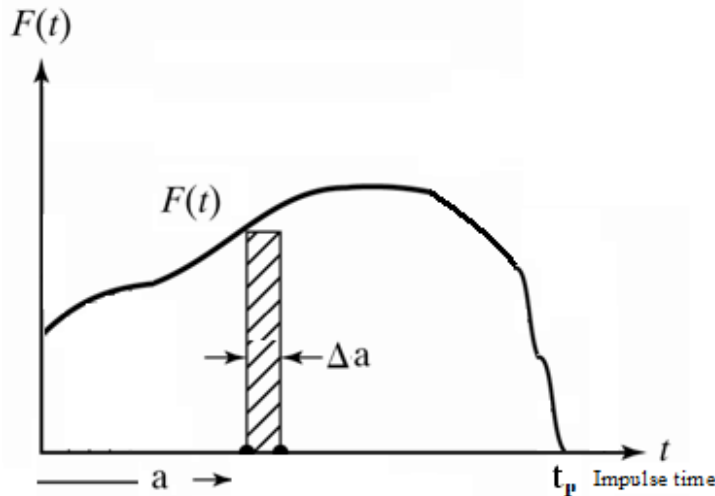
- Examine one of the impulses

- Cross-hatched in the figure
- At $t=a$ its strength is (the cross-hatched area):

$$\hat{F} = F(a) \Delta a$$

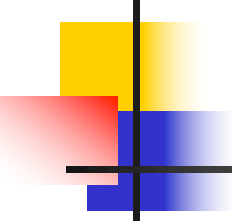
- The system's response at any time t , where $t > a$, is dependent upon the elapsed time $(t - a)$:

$$F(a) \Delta a g(t - a)$$



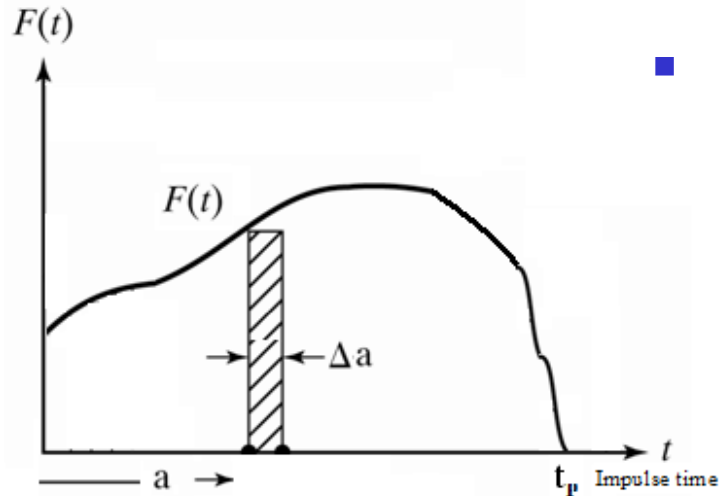
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Arbitrary Excitation



- The fundamental idea:
 - Consider the arbitrary excitation $F(t)$ to be just a train of successive impulsive excitations...
 - ... then apply the principle of superposition to find the response as the sum of the responses to this train of impulsive excitations

- Now sum up all the responses, for every single value of a ...
- This is where the integral part comes into play:



- Examine one of the impulses
 - Cross-hatched in the figure
 - At $t=a$ its strength is (the cross-hatched area):

$$\hat{F} = F(a) \Delta a$$

- The response at a time t is then obtained by combining all these contributions, for each $t > a$:

$$x_{Cnv}(t) = \sum F(a) g(t - a) \Delta a = \int_0^t F(a) g(t - a) da$$

[Cntd.]

Arbitrary Excitation

- Express in standard form:

$$x_{Cnv}(t) = \int_0^t F(\tau)g(t - \tau)d\tau$$

- Integral above: the “convolution integral” or “Duhamel’s integral”
- Another equivalent form in which it’s known (do change of variables...):

$$x_{Cnv}(t) = \int_0^t F(t - \lambda)g(\lambda)d\lambda$$

- Because τ and λ are dummy integration variables, I can say that

$$x_{Cnv}(t) = \int_0^t F(\tau)g(t - \tau)d\tau = \int_0^t F(t - \tau)g(\tau)d\tau$$

[Concluding Remarks]

Arbitrary Excitation

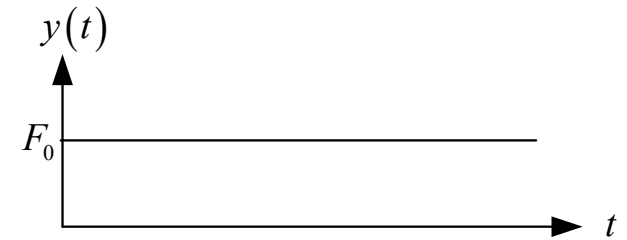
- You have a system and you are in a position to tell how it responds to a Dirac impulse (this response has been denoted by $g(t)$)
- You have some arbitrary excitation $F(t)$ applied to the system
- What have we just accomplished?
 - We can find the response $x_{Cnv}(t)$ of the system to the arbitrary excitation $F(t)$:

$$x_{Cnv}(t) = \int_0^t F(\tau)g(t - \tau)d\tau = \int_0^t F(t - \tau)g(\tau)d\tau$$

- In practical applications, when you have to evaluate one of these two integrals, shift the simpler of the two functions
 - See Examples...

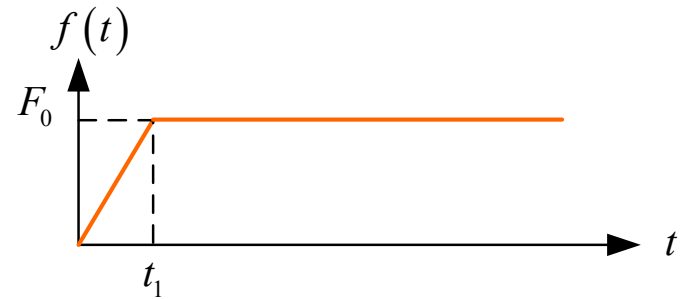


Example [AO]



- Determine the response of a 1DOF system to the step function
 - Work with an undamped system

Example [AO]



- Determine undamped response for a step function with rise time t_1
- Use a trick to break excitation function into two parts:

