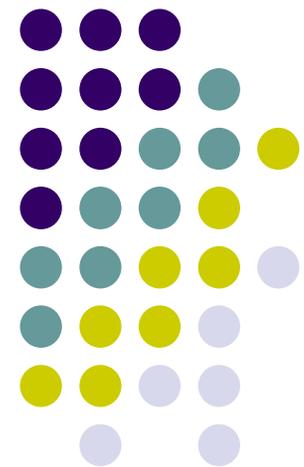


# ME751

## Advanced Computational Multibody Dynamics

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Solving Initial Value Problems  
Basic Concepts  
March 23, 2010



# Before we get started...



- Last Time:
  - Super briefly talked about the EOM when using Euler Angles
  - Discussed two classes of forces likely to be encountered in Engineering Apps
  - Inverse Dynamics Analysis
  - Equilibrium Analysis
- Today:
  - Discuss about the solution of differential equations
- Final Project:
  - Attach the one pager proposal to your Th HW
- I'll need a head count in the next two weeks regarding people who plan to make the trip to John Deere and NADS

# Dynamics Analysis



- Dynamics Analysis, Framework:
  - The state of a mechanical system (position, velocity) changes in time under the influence of internal and external **forces and/or prescribed motions**
  - The goal is to determine how the state of the system changes in time
  - Almost always you will only be able to determine the state of the mechanical system at a collection of grid points in time
    - That is, not everywhere, yet you can have as many grid points you wish (and afford)
  - Time evolution is obtained as the solution of the EOM (Newton-Euler equations derived before)

# Dynamics vs. Kinematics



- Kinematics Analysis
  - Prescribed motions exclusively determine how the system changes in time
    - The concept of force/torque does not factor in anywhere
  - For a Kinematics Analysis to be possible, the NDOF should be zero
  - Its solution provided at each time step by a sequence of 3 **algebraic** problems:
    - Nonlinear system of equations provides the position at each time step
    - Linear system of equations provides the velocity configuration at each time step
    - Linear system of equations provides the acceleration configuration at each time step
- Dynamics Analysis
  - External forces/torques dictate how the system evolves in time
  - It is more general than Kinematics:
    - A Kinematics problem can be solved using the methods of Dynamics, but not the other way around
  - Its solution obtained at each time step by numerical integration (solving a differential equation)

# Dealing With Dynamic Systems



- What have we done in ME751 so far?
  - We posed a dynamics problem
    - Discussed the elements of a mechanical system
    - Understood what their presence entails
    - Formulated a set of second order differential equations that governs the time evolution of the system
  - What's left: solving the differential equation to obtain the dynamics of the mechanical system
    - Four lectures: discuss solution techniques that produce an approximation of the solution of the differential problem

# 30,000 Feet Perspective



- When carrying out Dynamics Analysis, what you can compute is the acceleration of each part in the model
- Acceleration represents the second time derivative of your coordinates
- Somewhat oversimplifying the problem, in ME751 you get the second time derivative

$$\ddot{\mathbf{q}} = \mathbf{f}(\mathbf{q}, \dot{\mathbf{q}}, t)$$

- This represents a second order differential equation since it has two time derivatives taken on  $\mathbf{q}$
- Problem is reduced to a set of first order differential equations by introducing a helper variable  $\mathbf{v}$  (the velocity):

$$\dot{\mathbf{q}} = \mathbf{v}$$

- With this, the original second order differential problem becomes:

$$\dot{\mathbf{y}} = \mathbf{F}(\mathbf{y}, t) \quad \text{where} \quad \mathbf{y} = \begin{bmatrix} \mathbf{q} \\ \mathbf{v} \end{bmatrix} \quad \text{and} \quad \mathbf{F}(\mathbf{y}, t) \equiv \begin{bmatrix} \mathbf{v} \\ \mathbf{f}(\mathbf{v}, \mathbf{q}, t) \end{bmatrix}$$

# Numerical Integration

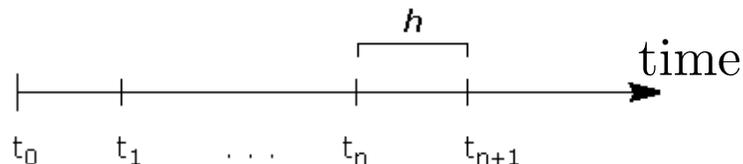
## ~The Problem~



- Initial Value Problem:  
(IVP)

$$\begin{cases} \dot{y}(t) = f(t, y) \\ y(t_0) = y_0 \end{cases}$$

- IVP: Stating the Problem
  - You are looking for a function  $y(t)$  that depends on time (changes in time), whose time derivative is equal to a function  $f(t, y)$  that is given to you (see equation above)
  - In other words, you are given the derivative of a function, can you tell what the function is?
- In ME751, the best you can hope for is to find an approximation of the unknown function  $y(t)$  at a sequence of discrete points (as many of them as you wish)
  - The numerical algorithm produces an approximation of the value of the unknown function  $y(t)$  at the each grid point. That is, the numerical algorithm produces an approximation for  $y(t_1)$ ,  $y(t_2)$ ,  $y(t_3)$ , etc.; i.e.,  $y_1$ ,  $y_2$ ,  $y_3$ , etc.



# Road Map (Tu & Th)



- Difference between ODE and IVP
- Basic Concepts in Numerical Integration
- Basic Methods for Numerical Integration
  - Runge-Kutta
  - AB & AM Methods
  - BDF Methods
- Text used:
  - Computer Methods for Ordinary Differential Equations and Differential-Algebraic Equations, by U. Ascher and L. Petzold, SIAM, 1998
  - On reserve at Wendt Library

# ODE vs. IVP

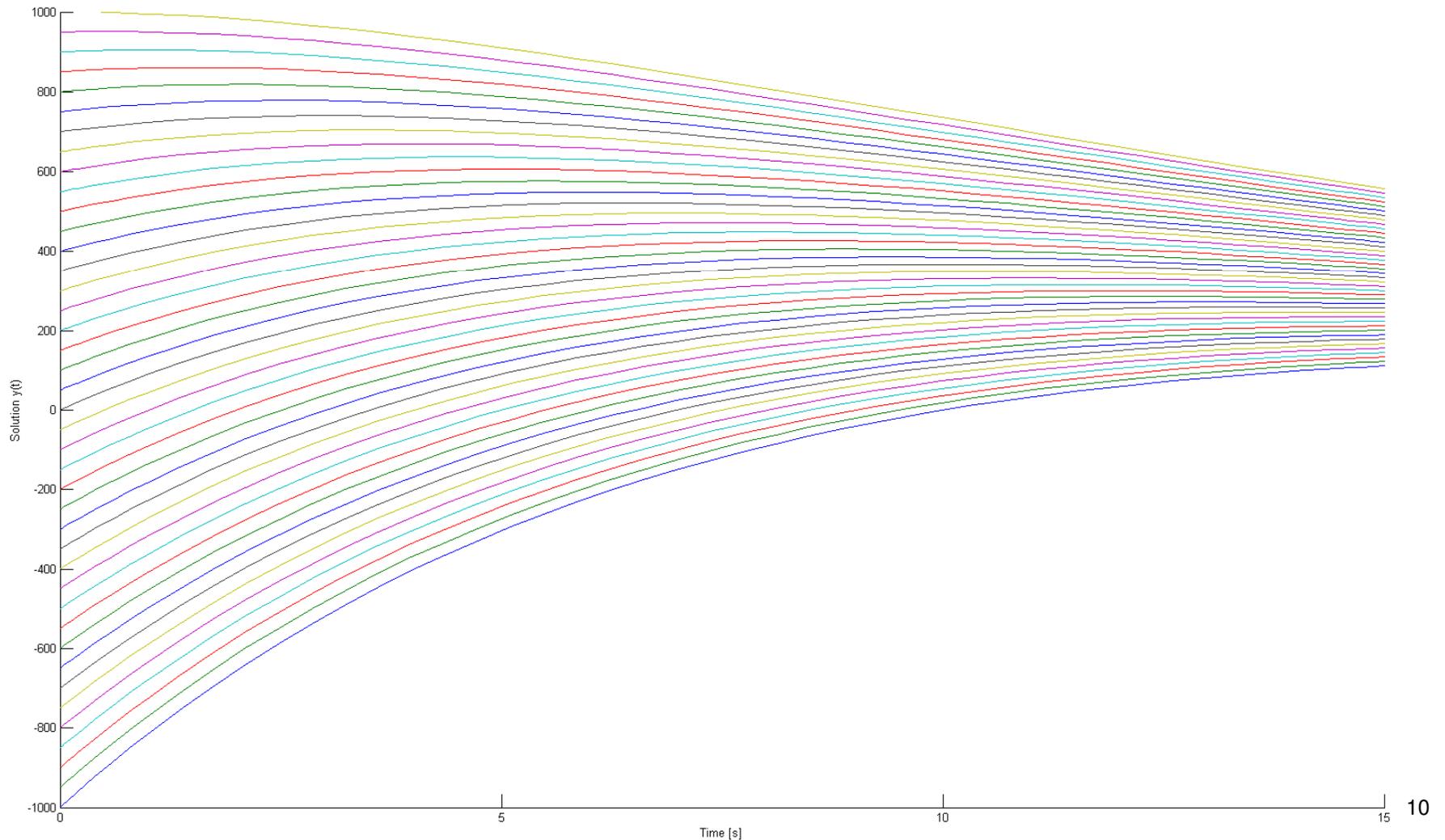


- Difference between ODE and IVP - Often a source of confusion
- Ordinary Differential Equation (ODE)
  - Typically, has an infinite number of solutions
- Initial Value Problem (IVP)
  - Is an ODE **plus** an initial condition (IC):
    - The IC: The unknown function assumes at time  $T=0$  a certain prescribed value
  - The solution for the IVP's that we'll deal with is **UNIQUE**
    - We'll assume that  $f(t,y)$  is well behaved (Lipschitz continuous)

# ODE: Infinite Number of Solutions



- ODE Problem:  $\dot{y}(t) = -0.1y(t) + 100e^{-0.1t}$
- A range of Initial Conditions (ICs) is specified:  $y_0 = [-1000 : 50 : 1000]$



# ODE vs. IVP

[Cntd.]



- Remember:
  1. IVP = ODE + IC
  2. IVP has a **UNIQUE** solution (unlike on ODE)
- Why is observation above important?
  - In general, when we start working on a numerical solution to a problem we better know that the problem we are trying to solve is well posed (has a solution and it is unique)
- Turns out that in the Dynamics Analysis we are dealing with a well posed problem (an IVP)
  - Focus then on finding a way to approximate its solution by using the computer
    - The computer will produce numbers that at each node of the time grid will approximate the value of the generalized coordinates and their velocity

# Initial Value Problems: Basic Concepts



- Truncation Error
- Accuracy
- Convergence
- 0-stability
- Order of a method
- Local Error
- Stability
  - Absolute stability
  - A-stable Integration Methods
  - L-Stable Integration Methods (Methods with Stiff Decay)

# Framework



- Interested in finding a function  $y(t)$  over an interval  $[0, b]$
- This function must satisfy the following IVP:

$$\begin{cases} \dot{y} &= f(t, y) \\ y(0) &= c \end{cases} \quad t \in [0, b]$$

- We assume that  $f$  is bounded and smooth, so that  $y$  exists, is unique, and smooth
- Given to you:
  - The constants  $c$  and  $b$
  - The function  $f(t, y)$ .

# Framework [Cntd.]

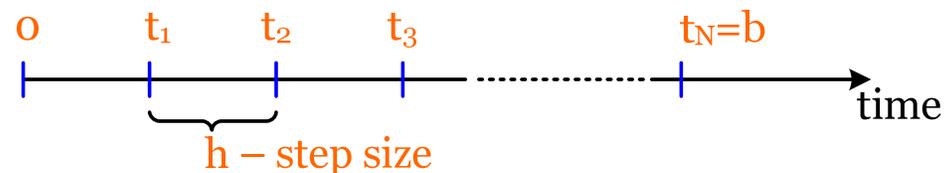


- The expression of the given function  $f(t,y)$  is typically far from being simple
  - Case 1:  $f(t,y)$  is simple, in very rare cases you can find  $y(t)$  analytically; i.e., find the exact solution of the IVP (the pen and paper case)
  - Case 2:  $f(t,y)$  is complex but nonetheless you have direct access to it.
    - Producing an exact solution is not possible, resort to numerical integration, solution approximated using the computer
  - Case 3:  $f(t,y)$  is so complex that you don't even have an expression for it. Instead you have another application (program) evaluate  $f(t,y)$  based on the value of  $t$  and  $y$  that you provide
    - Producing an exact solution is not possible, resort to numerical integration, solution approximated using several computer programs
    - This is what happens typically in multi-discipline engineering analysis

# Framework [Cntd.]



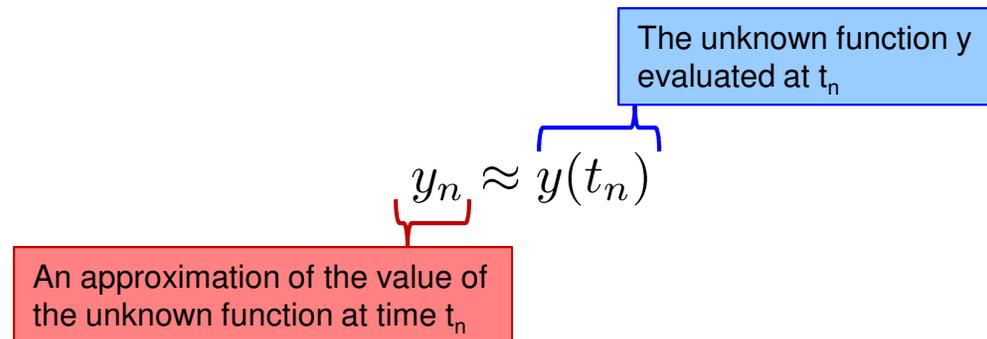
- Implications (in relation to the complexity of  $f(t,y)$ ):
  - Most likely you can't find the function  $y(t)$  analytically (using pen and paper)
  - You'll have to use a numerical approximation method to come up with approximations of the unknown function  $y(t)$ 
    - What you'll produce is an approximation of the value of  $y(t)$  at  $0, t_1, t_2$ , etc.
    - You are working on a grid of  $N$  time points, starting at  $0$  and ending at  $b$
    - For simplicity, we assume that the grid points are equally spaced by a value  $h$
    - This value is called the "integration step-size", usually denoted by  $h$  or  $\Delta t$





# Framework [Cntd.]

- **Remember this:** we'll approximate the value of  $y(t_n)$  by a value  $y_n$  that we'll learn how to obtain



- Notation: I'll use  $y^h$  to denote the set of values  $y_0, y_1, \dots, y_N$  that I come up with in my effort to approximate the unknown function  $y(t)$  at the grid points  $t_0, t_1, \dots, t_N$ :

$$y^h = \{y_0, y_1, \dots, y_N\}$$

- Quick Remark: a professional grade method for finding the approximate solution does not use an equally spaced grid of points  $0, t_1, t_2, t_3, \dots$ 
  - We'll stick with this assumption of equally spaced points though

# Basic Concepts: Truncation Error

[Preliminaries]



- We have our standard IVP: 
$$\begin{cases} \dot{y} & = & f(t, y) \\ y(0) & = & c \end{cases}$$

- Before talking about Truncation Error, let's introduce the simplest numerical integration scheme: Forward Euler
- Invoke a Taylor series expansion of function  $y(t)$  around  $t_{n-1}$ :

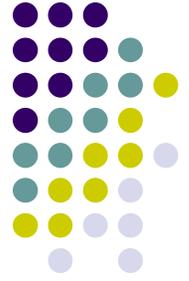
$$y(t_n) = y(t_{n-1}) + h\dot{y}(t_{n-1}) + \frac{1}{2}h^2\ddot{y}(t_{n-1}) + \dots$$

- Drop terms with powers of  $h$  of order 2 and higher
- Note that I won't use  $y(t_n)$  anymore, but I will now use  $y_n$ 
  - Make sure you understand this distinction between the exact value and approximate value...

[Basic Concepts]

# Truncation Error

[Preliminaries]



- I had this:

$$y(t_n) = y(t_{n-1}) + h\dot{y}(t_{n-1}) + \frac{1}{2}h^2\ddot{y}(t_{n-1}) + \dots$$

- But compute my solution like this

$$y_n = y_{n-1} + h\dot{y}_{n-1}$$

Truncation error

- Or equivalently,

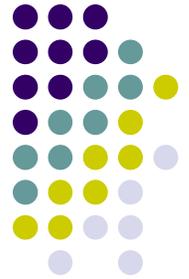
$$y_n = y_{n-1} + h f(t_{n-1}, y_{n-1})$$

- It's fair to expect that there is a slight difference between the exact value and numerical approximation:

## [Basic Concepts]

# Truncation Error

## [Definition]



- Consider how the solution is obtained:

$$\frac{y_n - y_{n-1}}{h} - f(t_{n-1}, y_{n-1}) = 0$$

- Note that in general, if you stick the actual solution in the equation above it is not going to be satisfied:

$$\frac{y(t_n) - y(t_{n-1})}{h} - f(t_{n-1}, y(t_{n-1})) \neq 0$$

- By definition, the quantity above is called the truncation error and is denoted by

$$\mathcal{N}(y, t_n, h) = \frac{y(t_n) - y(t_{n-1})}{h} - f(t_{n-1}, y(t_{n-1}))$$

- Note that this depends on the function ( $y$ ), the point where you care to evaluate the truncation error ( $t_n$ ), and the step size used ( $h$ )

[Basic Concepts]

# Accuracy Order



- Important remark:
  - Note that the truncation error depends on the integration scheme used to compute  $y^h$ :

$$\mathcal{N}(y, t_n, h)_{F.Euler} \neq \mathcal{N}(y, t_n, h)_{Runge-Kutta}$$

- Definition: an integration scheme is said to be accurate (or consistent) of order  $p$  for a positive  $p$  if

$$\mathcal{N}(y, t_n, h) = \mathcal{O}(h^p)$$

- By the way, a quantity “ $d$ ” is order  $h^p$  if there is a constant  $C$  so that when  $h$  is small enough, one has that

$$|d| \leq C h^p$$

[Basic Concepts]

# Accuracy Order



- Exercise:
  - Prove that Forward Euler is order 1 accurate, that is,

$$\mathcal{N}(y, t_n, h)_{F. Euler} = \mathcal{O}(h)$$

[Basic Concepts]

# Convergence



- Important question:
  - Does the numerical solution  $y^h$  actually converge to the unique solution of the IVP?
  - “Converge” means as  $h$  gets smaller and smaller, do you see the numerical solution at each  $t_n$  approaching the exact solution?

- More formal way to frame the convergence concept

- Define

$$e_n = |y_n - y(t_n)|, \quad n = 1, 2, \dots, N$$

- Note that  $N \cdot h = b$  (as  $h$  goes to zero,  $N$  keeps growing...)

- The numerical integration is said to be convergent of order  $p$  if

$$e_n = \mathcal{O}(h^p), \quad n = 1, 2, \dots, N$$

# Zero-stability (0-stability)



- There is a relationship between Accuracy and Convergence
- To capture this relationship we need to introduce the concept of zero-stability
- Definition: the numerical discretization scheme is 0-stable if there are positive constants  $h_0$  and  $K$  such that for any solutions  $x^h$  and  $z^h$ , obtained for  $h < h_0$ , one has that

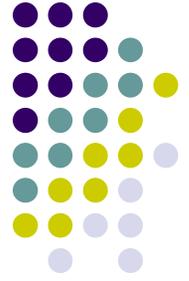
$$|x_n - z_n| \leq K \left( |x_0 - z_0| + \max_{1 \leq j \leq N} |\mathcal{N}(x, t_j, h) - \mathcal{N}(z, t_j, h)| \right), \quad 1 \leq n \leq N$$

# Exercise

- Prove that Forward Euler is a 0-stable discretization scheme



# Order “p” Convergence



- Theorem:

Consistency + 0-stability  $\Rightarrow$  Convergence

- Some more specifics:

- If the method is accurate of order  $p$  and 0-stable, then it is convergent of order  $p$ :

$$e_n = \mathcal{O}(h^p), \quad n = 1, 2, \dots, N$$

[Basic Concept]

# Local Error



- Imagine it as being the error that is noticed in \*one\* integration step
- Specifically, you start at time step  $t_{n-1}$ , from the point  $y_{n-1}$ 
  - The numerical solution finds for  $t_n$  the approximate value  $y_n$  (no surprise here)
  - The analytical solution that passes through the initial value  $y_{n-1}$  produces at  $t_n$  the value  $\bar{y}_n$  :

$$\begin{cases} \dot{\bar{y}} &= f(t, \bar{y}) \\ \bar{y}(t_{n-1}) &= y_{n-1} \end{cases}$$

- The local error is the difference

$$l_n = \bar{y}(t_n) - y_n$$

- Note the relationship that exists between the local error and local truncation error:

$$h \cdot |\mathcal{N}(\bar{y}, t_n, h)| = |l_n| \cdot (1 + \mathcal{O}(h))$$

[Basic Concept]

# Absolute Stability



- Convergence is good, but it tells me what happens if I work with smaller and smaller step-sizes  $h$
- In reality, I want to operate with values of  $h$  that are large
  - Recall that if  $h$  is small you have to take a very large number of integration steps to cover the interval  $[0,b]$
- The relevant question: How large can I consider  $h$  yet know for a fact that I don't get garbage approximations of the solution?
  - The answer to this question is posed to each numerical discretization scheme
  - Moreover, it is posed in conjunction with a test problem:

$$\begin{cases} \dot{y} &= \lambda y \\ y(0) &= 1 \end{cases}$$

[Basic Concept, Short Detour]

# Absolute Stability



- Example: mass-spring-damper oscillation

$$\begin{aligned}\ddot{x} + 2\zeta\omega_n\dot{x} + \omega_n^2x &= 0 \\ x(0) &= x_0 \\ \dot{x}(0) &= \dot{x}_0\end{aligned}$$

- Case 1: underdamped system

$$\lambda_1 = -\omega_n\zeta - j\omega_d \quad \lambda_2 = -\omega_n\zeta + j\omega_d$$

- Case 2: overdamped system

$$\lambda_1 = -\omega_n\zeta - \omega_n\sqrt{\zeta^2 - 1} \quad \lambda_2 = -\omega_n\zeta + \omega_n\sqrt{\zeta^2 - 1}$$

- Case 3: critically damped system

$$\lambda_1 = \lambda_2 = -\omega_n$$

[Basic Concept]

# Absolute Stability



- Example, the concluding remark:
  - In mechanical engineering,  $\lambda$  assumes values for which typically

$$\Re(\lambda) < 0$$

- It better be that our numerical discretization scheme leads to numerical solutions that can handle the test IVP for negative values of  $\lambda$  (or its real part, when dealing with complex values...)
- It turns out that this is not trivial
  - Take Forward Euler for a spin to see how problems crop up...

## Example: ( $\lambda=-100$ )

$$\left. \begin{aligned} \dot{y} &= -100y \\ y(0) &= 1 \end{aligned} \right\} \Rightarrow y(t) = e^{-100t}$$



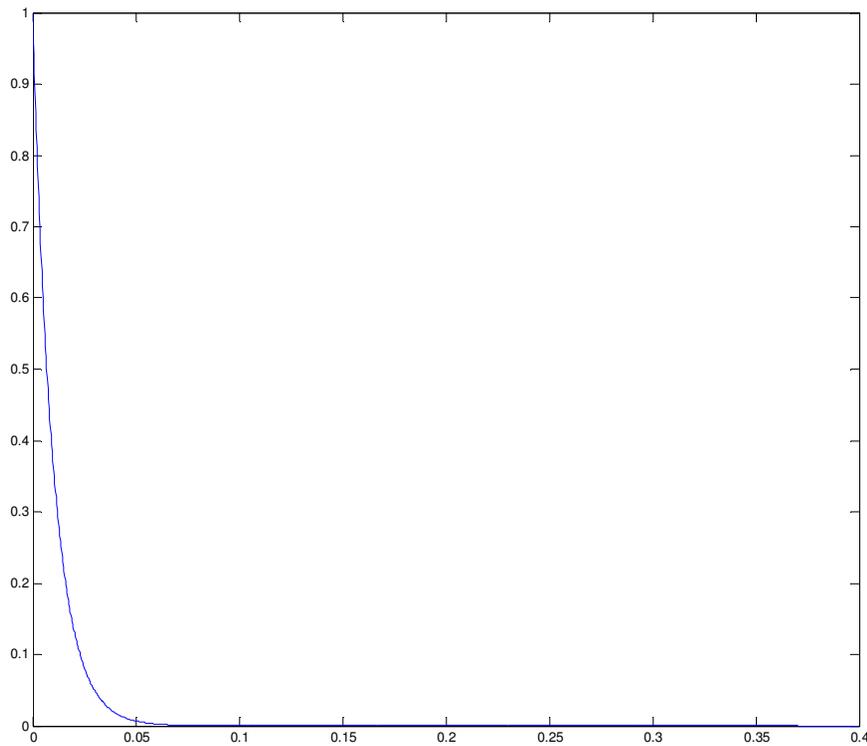
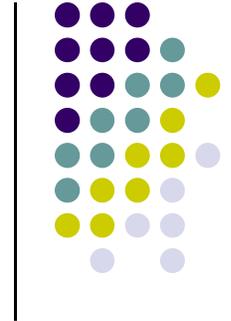
- Integrate 5 steps using forward Euler formula:  $\Delta t=0.002$ ,  $\Delta t=0.01$ ,  $\Delta t=0.03$
- Compare the errors between numerical and analytical solutions

$\Delta t=0.002$ : Error	$\Delta t=0.01$ : Error	$\Delta t=0.03$ : Error
0	0	0
0.01873075307798	0.36787944117144	2.04978706836786
0.03032004603564	0.13533528323661	-3.99752124782333
0.03681163609403	0.04978706836786	8.00012340980409
0.03972896411722	0.01831563888873	-15.99999385578765
0.04019944117144	0.00673794699909	32.00000030590232

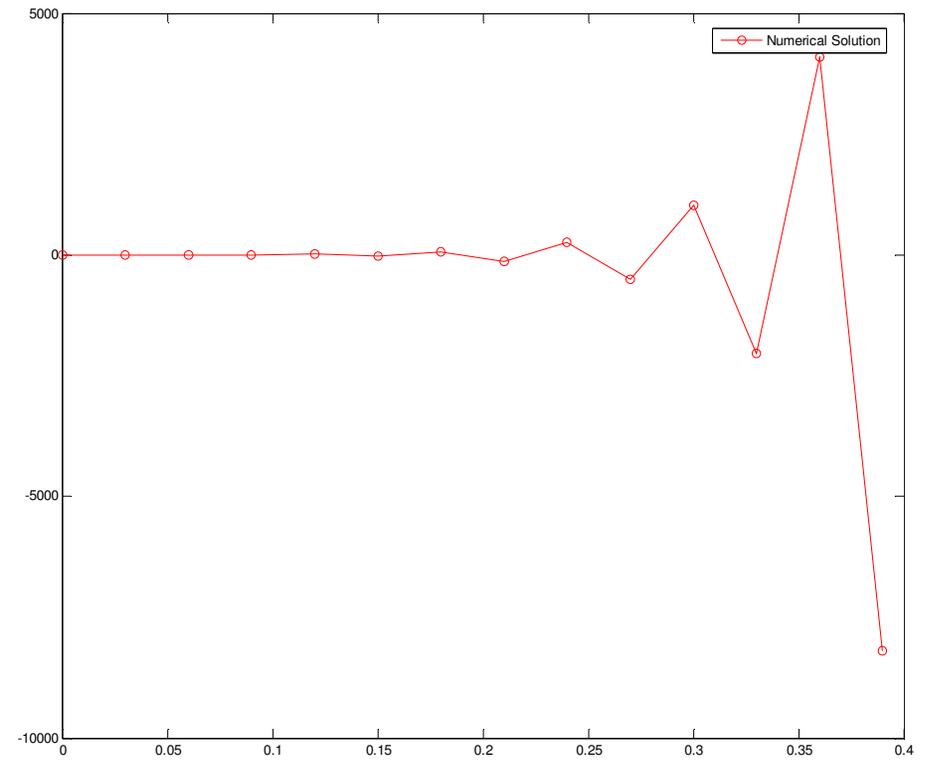


**Example:**  
( $\lambda=-100$ )

$$\left. \begin{aligned} \dot{y} &= -100y \\ y(0) &= 1 \end{aligned} \right\} \Rightarrow y(t) = e^{-100t}$$



Analytical Solution



Forward Euler

**( $\Delta t=0.03$ )**

# Basic Concept: Absolute Stability



- One quick way to see that things went south
- First note that
  - if  $\Re(\lambda) < 0$ , then  $0 < y(t_n) < y(t_{n-1}) < \dots < y(t_1) < y(t_0) = 1$

- You'd expect that the numerical solution will mirror this behavior:

$$0 < y_n < y_{n-1} < \dots < y_1 < y_0 = 1$$

- Use Forward Euler to express  $y_n$  as a function of  $y_{n-1}$  and require that the condition  $y_n < y_{n-1}$  to obtain that

$$|1 + h\lambda| < 1$$

# Basic Concept: Absolute Stability



- Recall what happened
  - We started with the test problem
  - We required that for the test problem, the numerical approximation should behave like the solution. That is, we required that

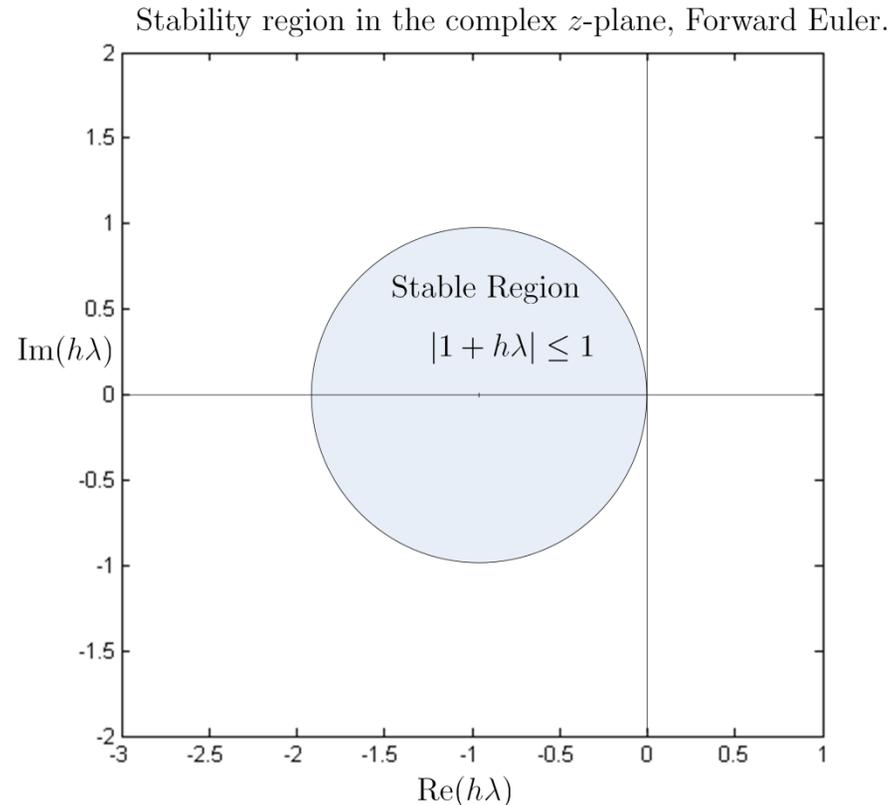
$$y_n \leq y_{n-1}$$

- We used the discretization scheme (Forward Euler, in our case) to express how  $y_n$  is related to  $y_{n-1}$
- This led to a condition that the step size should satisfy, specifically, for Forward Euler, we obtained that

$$|1 + h\lambda| < 1$$

- What we just did was to determine the region of absolute stability of Forward Euler
- The KEY point: to get the absolute stability region for other integration methods, you have to use that integration method to express how  $y_n$  is related to  $y_{n-1}$

# Basic Concept: Absolute Stability



- The step size  $h$  should be such that  $h\lambda$  lands into the shaded circle
- Note that a very negative value for  $\lambda$  will require a very small value of  $h$  so that the product  $h\lambda$  is inside the circle

# Accuracy vs. Stability



- Note that Forward Euler is accurate to order 1
  - That is, locally,

$$\mathcal{N}(y, t_n, h)_{F. Euler} = \mathcal{O}(h)$$

- This is an asymptotic and \*local\* result, which holds better as h gets smaller
- For the test IVP, the local error compounds due to the particular form of the problem that we work with (the test IVP)
- This compounding and the fact that h does not assume small values (you try to work w/ large h) leads to the phenomenon of loss of stability
- To conclude, there is no contradiction here (the numerical scheme being order 1 accurate yet losing stability for large values of h)

# Accuracy vs. Stability: The Concept of Stiffness



- Recall that the size/shape of the stability region (SR) is specific to each discretization scheme
- For some discretization schemes the SR is ridiculously small
  - Forward Euler is one of them
- A small SR to start with, combined with a problem for which  $\lambda$  is very negative leads to unreasonably small values of  $h$ 
  - Such a problem is called a “stiff” IVP
- In this case it is \*not\* the accuracy concerns that restrict the value of the step-size  $h$ , but rather the stability issue prevails
- Note that it can be the case that if you change the integration method, the stiff problem is just fine (if the stability region is generous)



# Example:

- Use Forward Euler to find an approximation of the solution of the following IVP:

$$\begin{cases} \dot{y} = -100y + \sin(t) \\ y(0) = 0 \end{cases} \quad t \in [0, 8]$$

# Example [Cntd.]

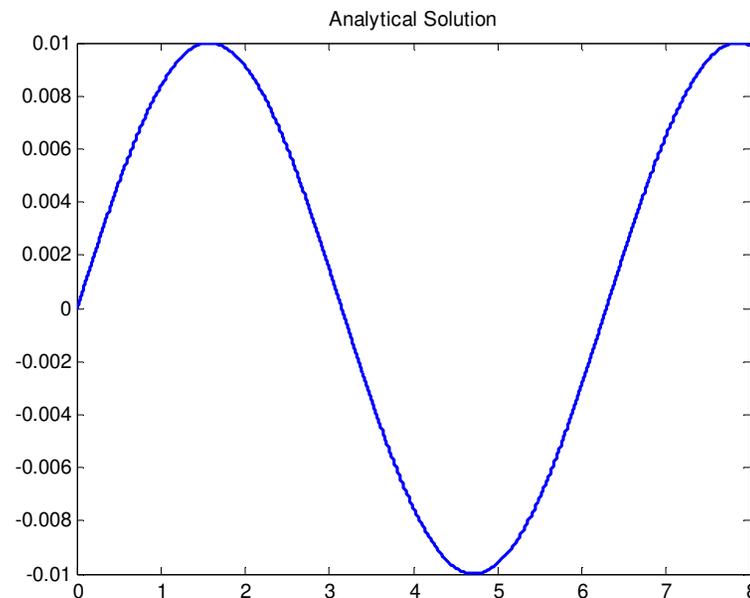


- Easy to figure out that the exact solution is

$$y(t) = \frac{1}{10001} (100 \sin(t) - \cos(t) + e^{-100t})$$

- For all purposes, the solution can be rewritten as

$$y(t) \approx 0.01 \sin(t - \phi_0) \quad \text{where} \quad \tan \phi_0 = \frac{1}{100}$$



# Example [Cntd.]



- Note that for the given IVP,  $\lambda = -100$ , which suggests we'll have to work with small step-sizes...
- Note in the plot that the contribution of the exponential is not even seen in the MATLAB plot
- You'd expect that since the exponential component of the solution goes away so quickly one could use Forward Euler and have no difficulties, which is not true... (see next slide)

```
% Solves the IVP y_dot = -100*y + sin(t) and y(0)=0
% yExact is the analytical solution
% yFE is the solution obtained with Forward Euler
% yBE is the solution obtained with Backward Euler
%
% Input: the integration step-size

h = input('Input step size h:');
tend = 8;
tm=0:h:tend;

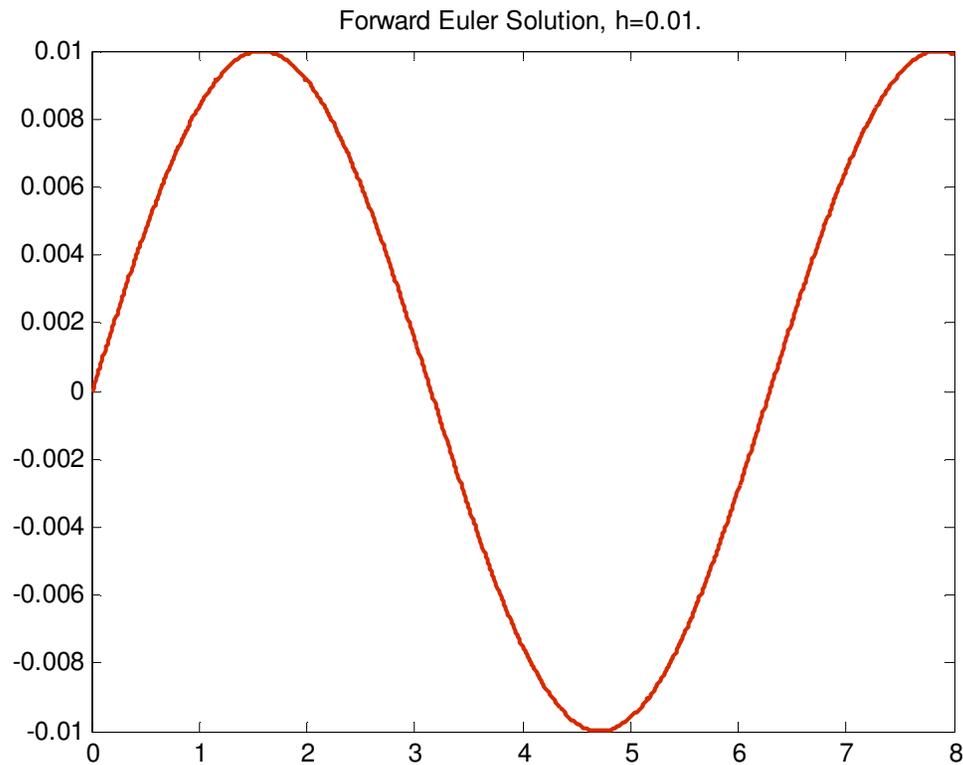
yExact = 1/10001*(100*sin(tm)-cos(tm)+exp(-100*tm));

yFE = zeros(size(tm));
yBE = zeros(size(tm));

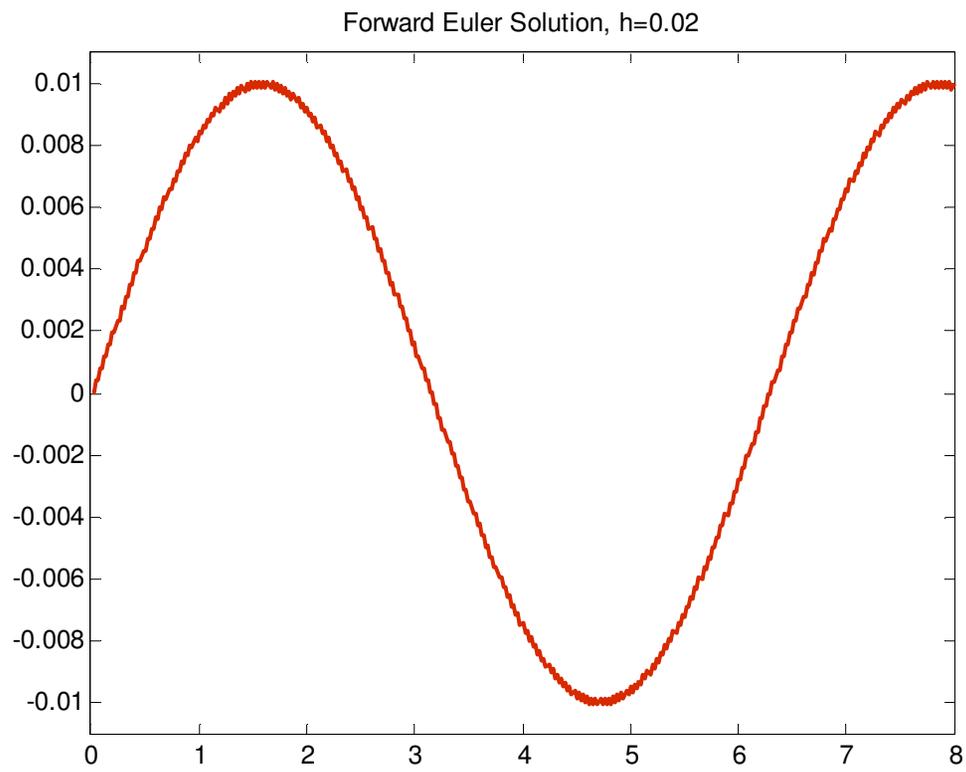
for i=2:1:length(yFE)
    yFE(i) = yFE(i-1)*(1-100*h) + h*sin(tm(i-1));
end

dummyINV = 1/(1+100*h);
for i=2:1:length(yBE)
    yBE(i) = yBE(i-1)*dummyINV + h*dummyINV*sin(tm(i));
end
```

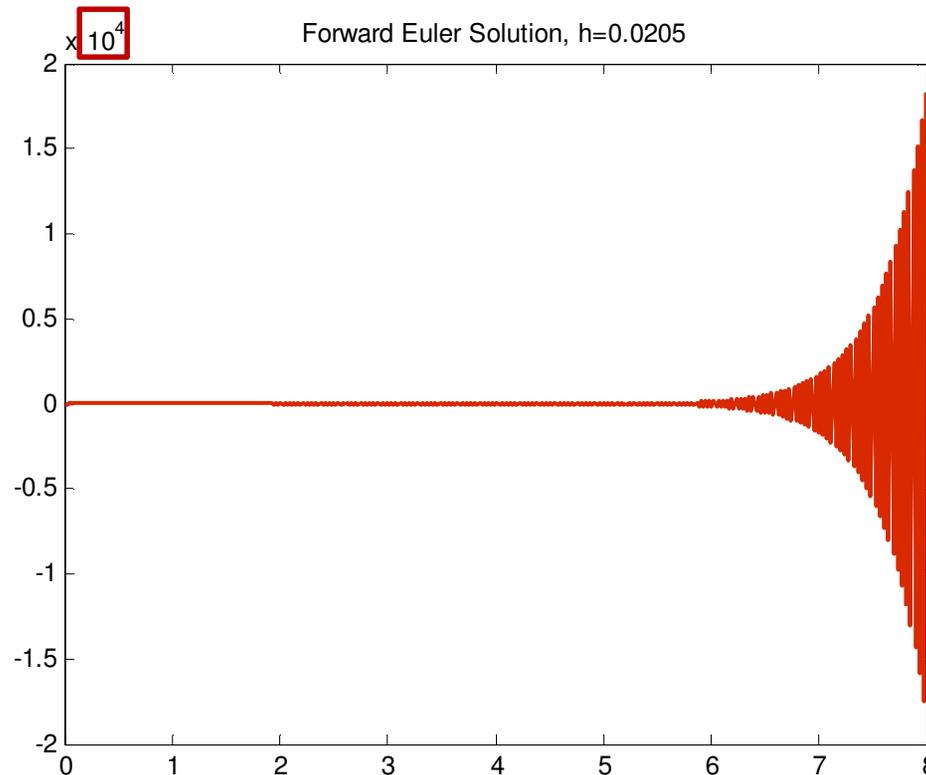
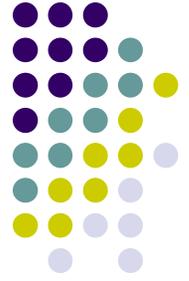
# Example, Approached with Forward Euler: $h=0.01$ s



# Example, Approached with Forward Euler : $h=0.02$ s



# Example, Approached with Forward Euler: $h=0.0205$ s



- As soon as you go beyond the limit value  $h=0.02$  (that goes hand in hand for Forward Euler with  $\lambda=-100$ ), you run into trouble
- Note that this happens even though the contribution of the exponential goes away very fast...

# Example, Approached with Forward Euler



- Conclusion
  - For this type of problem with very negative  $\lambda$ , Forward Euler is bad
    - The step size is significantly limited on stability grounds
- Qualitative definition:
  - An IVP where Forward Euler behaves bad is called STIFF IVP